

Eisenstein series attached to small automorphic representations

Henrik Gustafsson

Automorphic forms, mock modular forms and string theory
Simons Center for Geometry and Physics 2016

Based on

Small automorphic representations and degenerate Whittaker vectors

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[arXiv:1412.5625](https://arxiv.org/abs/1412.5625) [math.NT]

[GKP14]

Journal of Number Theory 166 (Sep, 2016) 344–399

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$SL(n)$

E_6, E_7, E_8

Outline

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- Why do string theorists study Eisenstein series?

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- Why do *string theorists* study Eisenstein series?
- What *Fourier coefficients* are we interested in and why?
- How can we compute them?
- What happens for *small automorphic representations*?
- What's next?

Motivation

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- Hecke eigenvalues
- Point counts of elliptic curves
- Langlands program
L-functions | The Langlands–Shahidi method

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Scattering amplitudes | Black hole microstate counting
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String theory



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String theory



World-sheet
 Σ

String theory



World-sheet
 Σ

Typical string length: l_s



String theory



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 Σ

Typical string length: l_s



$$\alpha' = l_s^2$$

String theory

Space-time is described by a Riemannian manifold M

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String theory = dynamics of the embedding maps

$$X : \Sigma \rightarrow M$$

world-sheet space-time

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Consistency requires: 10-dimensional M

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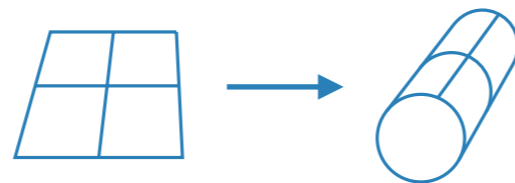
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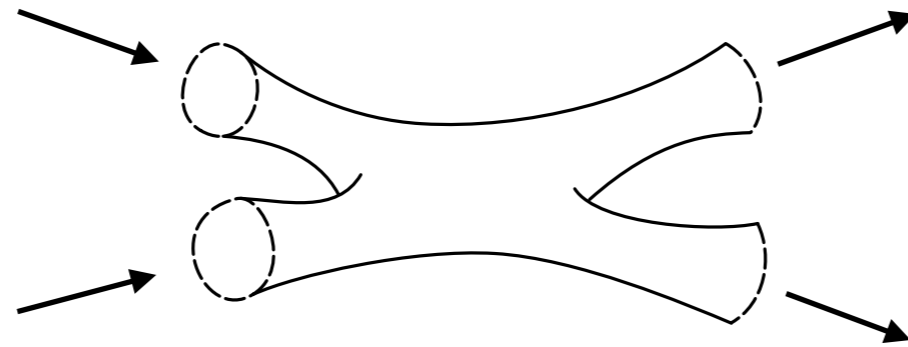
Consistency requires: 10-dimensional M

Toroidal compactifications

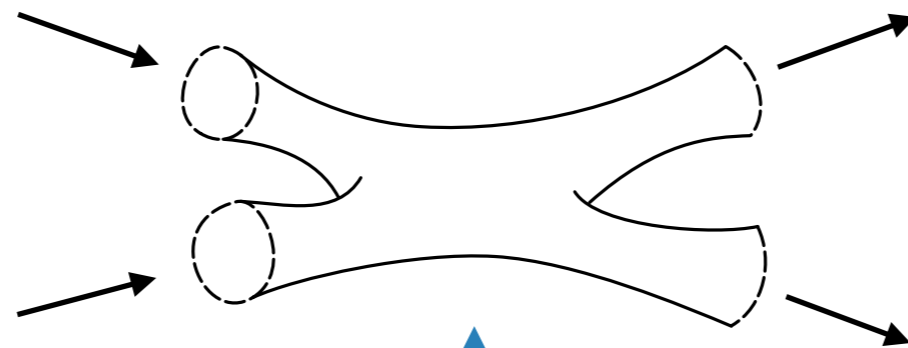


D dimensions

Interactions

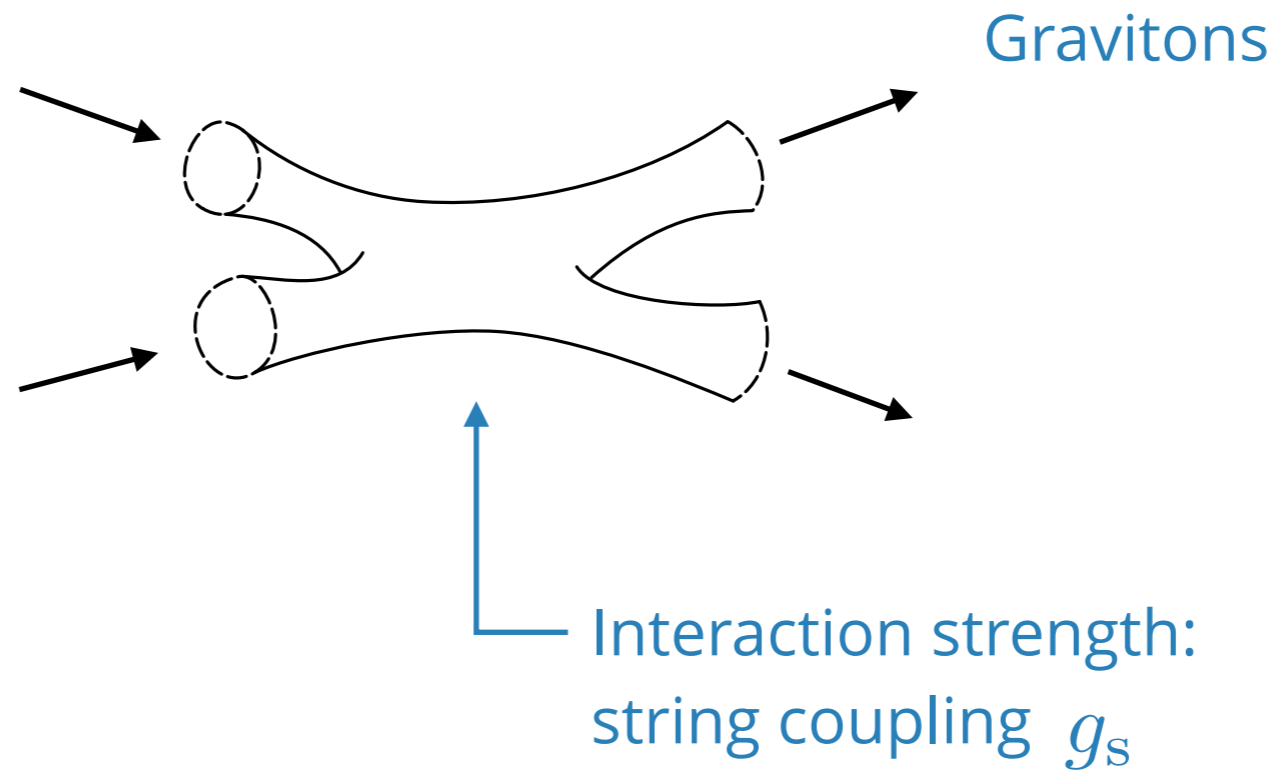


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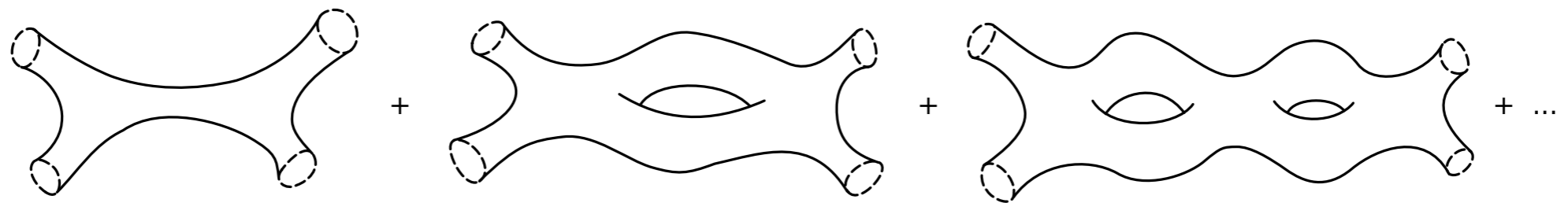


Interaction strength:
string coupling g_s

Interactions

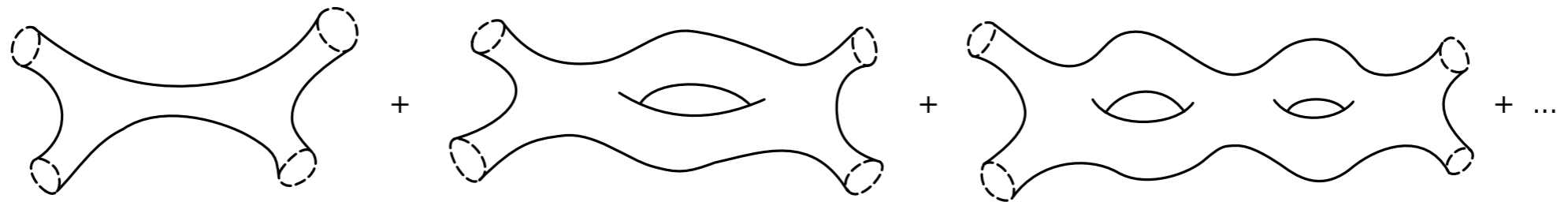


Interactions



Interactions

Weighted by: $g_s^{-\chi_E}$ $-\chi_E = 2(\text{genus} - 1) + \text{boundaries}$

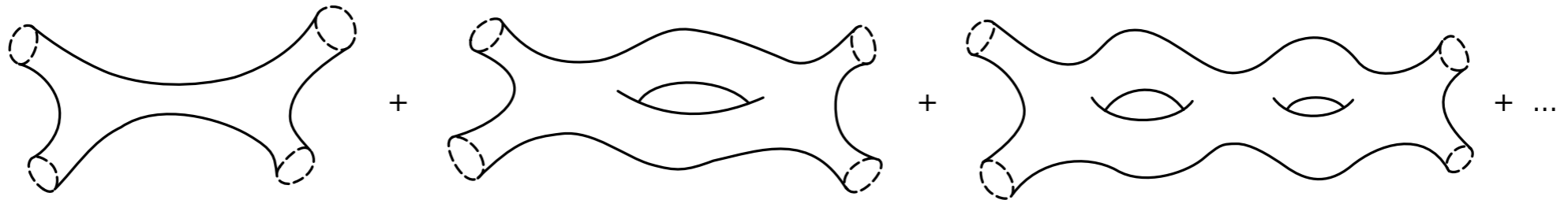


Interactions

Weighted by: $g_s^{-\chi_E}$

Euler characteristic

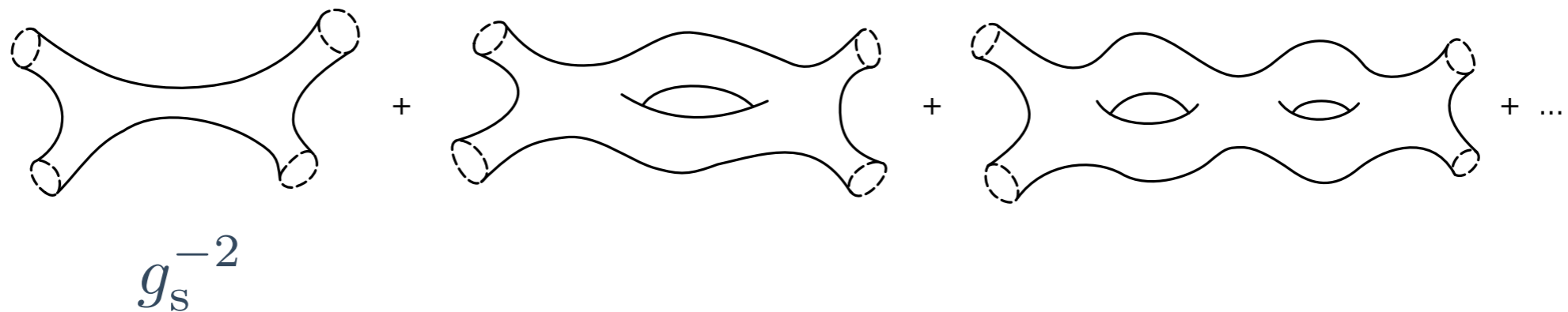
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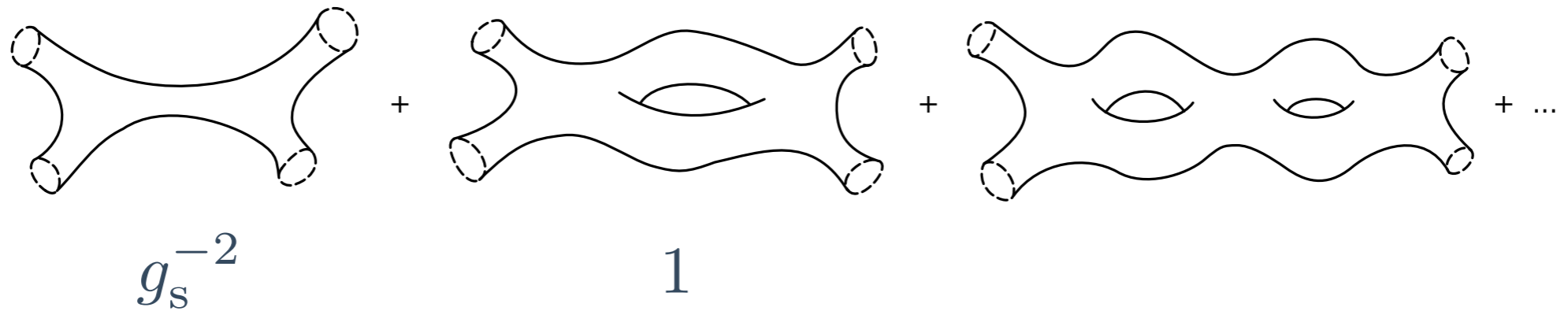


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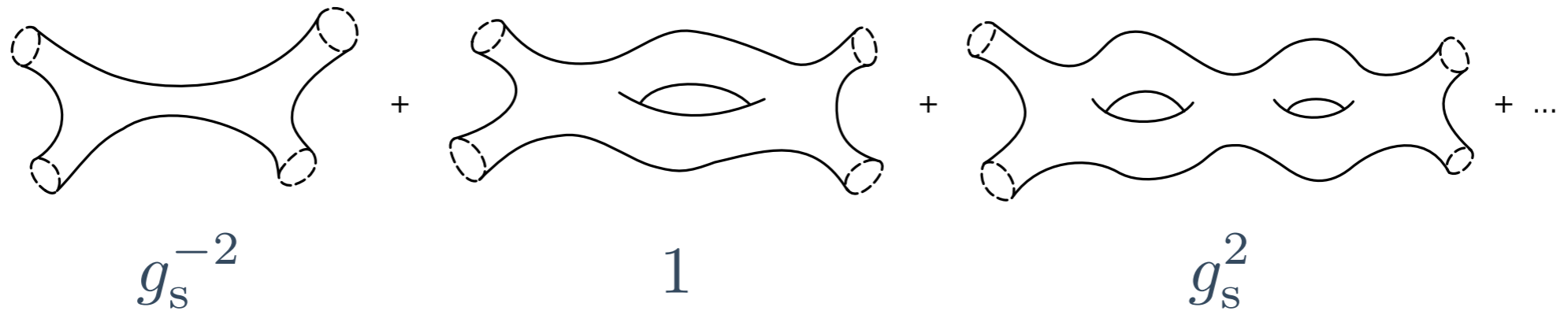
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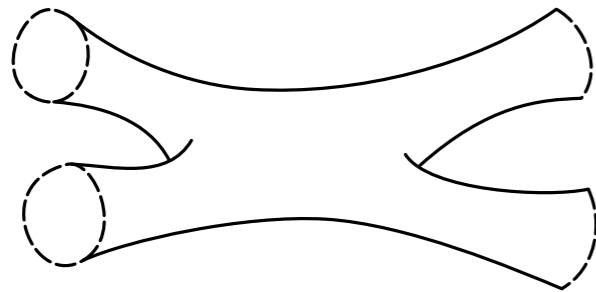
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Interactions

Gravitons in D dimensions



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$$R + (\alpha')^3 \mathcal{E}_{(0,0)}^{(D)}(g) R^4 + (\alpha')^5 \mathcal{E}_{(1,0)}^{(D)}(g) D^4 R^4 + (\alpha')^6 \mathcal{E}_{(0,1)}^{(D)}(g) D^6 R^4 + \dots$$

Expansion
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Contractions of derivatives and 4 Riemann tensors (known)

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Moduli space

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8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$
7	$SL(5, \mathbb{R})$	$SO(5)$
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5	$E_6(\mathbb{R})$	$USp(8)/\mathbb{Z}_2$
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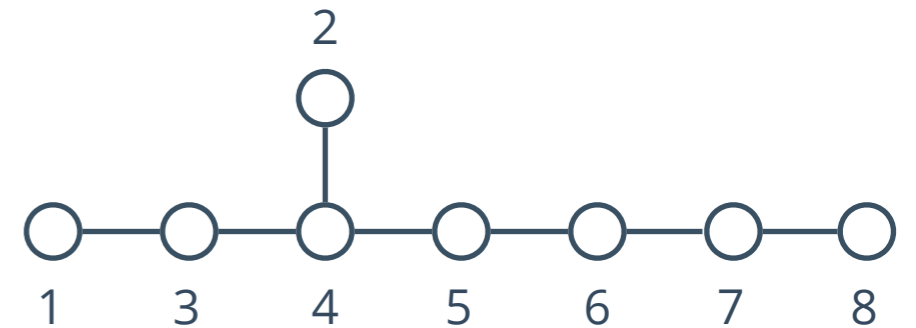
[Cremmer-Julia]

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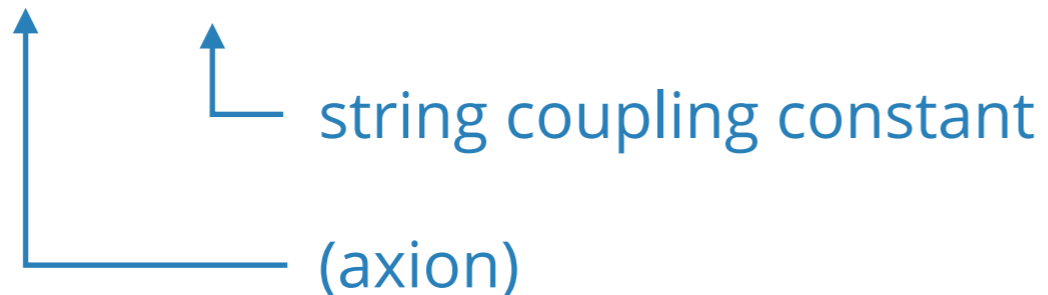
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 string coupling constant

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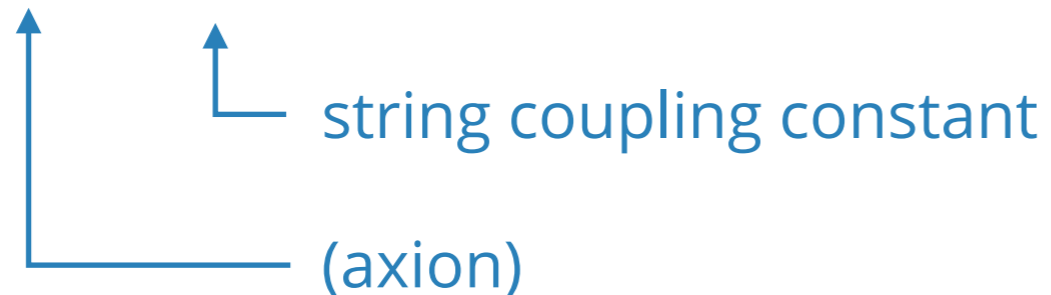
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$$\mathcal{E}_{(p,q)}(\tau) = \mathcal{E}_{(p,q)}^{(10)}(g)$$

U-duality

$G(\mathbb{R}) \curvearrowright \mathcal{M}_{\text{classical}}$ classical symmetry

[Hull-Townsend]

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Quantization of charges

[Hull-Townsend]

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All observables are invariant under $G(\mathbb{Z})$

[Hull-Townsend]

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$$\mathcal{E}_{(0,0)}^{(D)}(g), \mathcal{E}_{(1,0)}^{(D)}(g), \mathcal{E}_{(0,1)}^{(D)}(g) : G(\mathbb{Z}) \backslash G(\mathbb{R}) / K \rightarrow \mathbb{C}$$

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- (B) φ is an eigenfunction under right-translations of $k \in K$

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- (C) φ is an eigenfunction to all G -invariant differential operators

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- (C) Z-finiteness: $\dim(\text{span}\{X\varphi(g) \mid X \in \mathcal{Z}(\mathfrak{g}_{\mathbb{C}})\}) < \infty$

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- (C) Z-finiteness: $\dim(\text{span}\{X\varphi(g) \mid X \in \mathcal{Z}(\mathfrak{g}_{\mathbb{C}})\}) < \infty$
- (D) Growth: for any norm $\|\cdot\|$ on $G(\mathbb{R})$ there exists a positive integer n and constant C such that $|\varphi(g)| \leq C\|g\|^n$

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Supersymmetry constraints



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$$\left(\Delta - \frac{3}{4}\right) \mathcal{E}_{(0,0)}(\tau) = 0$$

[Green-Sethi]

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$$\left(\Delta - 12\right) \mathcal{E}_{(0,1)}(\tau) = -\left(\mathcal{E}_{(0,0)}(\tau)\right)^2 \quad \text{[Green-Vanhove]}$$

Supersymmetry constraints



10 dimensions: $\Delta = 4\tau_2^2 \frac{\partial}{\partial \tau} \frac{\partial}{\partial \bar{\tau}}$ Laplacian

(C) ✓ $(\Delta - \frac{3}{4})\mathcal{E}_{(0,0)}(\tau) = 0$ [Green-Sethi]

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Not an automorphic form in a strict sense

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Not an automorphic form in a strict sense

Similarly for lower dimensions

Eisenstein series

Eisenstein series

$$B(\mathbb{R}) = \left\{ \begin{pmatrix} * & * \\ 0 & * \end{pmatrix} \in SL(2, \mathbb{R}) \right\} \quad \text{Borel subgroup}$$

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Multiplicative character
trivially extended to $G(\mathbb{R})$
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$$E(s; \tau) = \sum_{\gamma \in B(\mathbb{Z}) \backslash SL(2, \mathbb{Z})} \chi(\gamma(\tau)) = \sum_{\gamma \in B(\mathbb{Z}) \backslash SL(2, \mathbb{Z})} \frac{\tau_2^s}{|c\tau + d|^{2s}}$$

$$\tau = \tau_1 + i\tau_2 \quad \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{Z}) \quad \gamma(\tau) = \frac{a\tau + b}{c\tau + d}$$

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$$E(s; \tau) = \tau_2^s + \frac{\xi(2s - 1)}{\xi(2s)} \tau_2^{1-s} + \frac{2\tau_2^{1/2}}{\xi(2s)} \sum_{m \neq 0} |m|^{s-1/2} \sigma_{1-2s}(m) K_{s-1/2}(2\pi |m| \tau_2) e^{2\pi i m \tau_1}$$

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Completed Riemann zeta function

$$\xi(s) = \pi^{-s/2} \Gamma(s/2) \zeta(s)$$

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Bessel function
of the second kind

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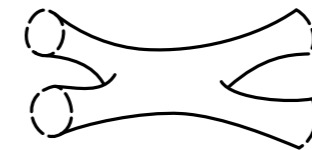
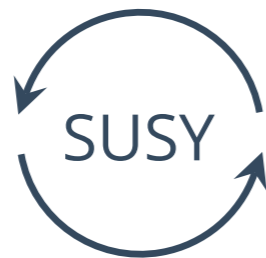
$$(\Delta - s(s - 1))E(s; \tau) = 0 \quad E(s; \tau) \sim \tau_2^s \quad g_s = \tau_2^{-1} \rightarrow 0$$

[Green-Gutperle, Pioline, Green-Russo-Vanhove]

Eisenstein series

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$$E(s; \tau) \sim \tau_2^s \quad g_s = \tau_2^{-1} \rightarrow 0$$



$$(\Delta - \frac{3}{4})\mathcal{E}_{(0,0)}(\tau) = 0$$

$$\mathcal{E}_{(0,0)}(\tau) \sim 2\zeta(3)\tau_2^{3/2}$$

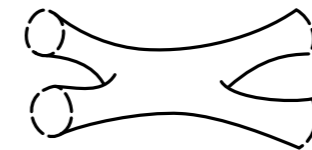
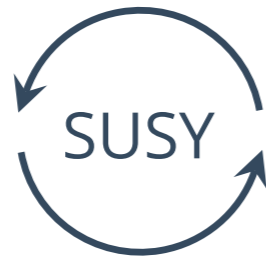
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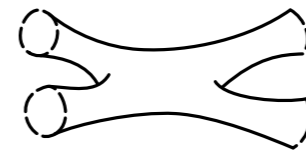
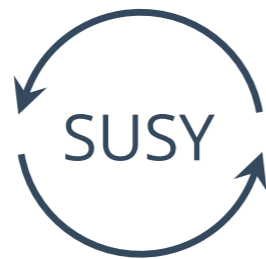
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$\mathcal{E}_{(0,1)}(\tau)$ as a sum over images $\sum_{B(\mathbb{Q}) \backslash G(\mathbb{Z})}$ but not of a character χ

[Green-Miller-Vanhove]

Extracting physical information

Expand Bessel function in g_s

$$\tau = \chi + ig_s^{-1}$$

Extracting physical information

Expand Bessel function in g_s $\tau = \chi + ig_s^{-1}$

 Interaction strength

Extracting physical information

Expand Bessel function in g_s $\tau = \chi + ig_s^{-1}$

The diagram illustrates the expansion of a Bessel function in terms of the interaction strength g_s . The text "Expand Bessel function in g_s " is followed by the equation $\tau = \chi + ig_s^{-1}$. A blue arrow points from the text "Expand Bessel function in g_s " to the symbol g_s . Another blue arrow points from the text "(axion)" to the term ig_s^{-1} . A third blue arrow points from the text "Interaction strength" to the symbol g_s .

Interaction strength

(axion)

Extracting physical information

Expand Bessel function in g_s $\tau = \chi + ig_s^{-1}$
Interaction strength (axion)

$$\mathcal{E}_{(0,0)}(\tau) = 2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-2\pi|m|g_s^{-1} + 2\pi im\chi} \left[1 + \mathcal{O}(g_s) \right]$$

Extracting physical information

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.....
Perturbative
(zero-mode)

[Green-Gutperle]

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Perturbative
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Non-perturbative
(remaining modes)

$$e^{-\frac{1}{g_s}}$$

[Green-Gutperle]

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Instanton action

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Non-perturbative
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Non-perturbative
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$$\sigma_s(m) = \sum_{d|m} d^s$$

Sums over the number of ways the charge m can be factorised into two integers

Extracting physical information

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Sums over the number of ways the charge m can be factorised into two integers



wrapping number and charge
of a T-dual D-particle

[Green-Gutperle]

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arithmetic information
 p -adic part

Sums over the number of ways the charge m can be factorised into two integers

wrapping number and charge
of a T-dual D-particle

[Green-Gutperle]

Lower dimensions

Lower dimensions

D	$G(\mathbb{R})$	K	$G(\mathbb{Z})$
10	$SL(2, \mathbb{R})$	$SO(2)$	$SL(2, \mathbb{Z})$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$	$SL(2, \mathbb{Z}) \times \mathbb{Z}_2$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$	$SL(3, \mathbb{Z}) \times SL(2, \mathbb{Z})$
7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5)) / \mathbb{Z}_2$	$Spin(5, 5; \mathbb{Z})$
5	$E_6(\mathbb{R})$	$USp(8) / \mathbb{Z}_2$	$E_6(\mathbb{Z})$
4	$E_7(\mathbb{R})$	$SU(8) / \mathbb{Z}_2$	$E_7(\mathbb{Z})$
3	$E_8(\mathbb{R})$	$Spin(16) / \mathbb{Z}_2$	$E_8(\mathbb{Z})$

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8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$	$SL(3, \mathbb{Z}) \times SL(2, \mathbb{Z})$
7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5)) / \mathbb{Z}_2$	$Spin(5, 5; \mathbb{Z})$
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3	$E_8(\mathbb{R})$	$Spin(16) / \mathbb{Z}_2$	$E_8(\mathbb{Z})$

$$E(\chi; g) = \sum_{\gamma \in B(\mathbb{Z}) \setminus G(\mathbb{Z})} \chi(\gamma g)$$

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U

Parabolic subgroups

Fourier expand
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Unipotent subgroup U



Choice of parabolic subgroup P

Parabolic subgroups

Fourier expand
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Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

$\langle \Sigma \rangle$ generated root system

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

$\langle \Sigma \rangle$ generated root system

$$\mathfrak{g}_\alpha = \{g \in \mathfrak{g} \mid [h, g] = \alpha(h)g \quad \forall h \in \mathfrak{h}\}$$

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

$\langle \Sigma \rangle$ generated root system

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Cartan subalgebra

Parabolic subgroups

Fourier expand
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Cartan subalgebra

$$\mathfrak{p} = \mathfrak{h} \oplus \bigoplus_{\alpha \in \Delta(\mathfrak{p})} \mathfrak{g}_\alpha$$

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

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$$\Delta(\mathfrak{p}) = \Delta_+ \cup \langle \Sigma \rangle$$

Parabolic subgroups

Fourier expand
in different directions



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Positive roots

Parabolic subgroups

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Levi decomposition

Parabolic subgroups

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Levi decomposition

$$\mathfrak{p} = \mathfrak{l} \oplus \mathfrak{u}$$

Parabolic subgroups

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Levi decomposition

$$\mathfrak{p} = \mathfrak{l} \oplus \mathfrak{u} \quad \mathfrak{l} = \mathfrak{h} \oplus \bigoplus_{\alpha \in \langle \Sigma \rangle} \mathfrak{g}_{\alpha} \quad \mathfrak{u} = \bigoplus_{\alpha \in \Delta(\mathfrak{u})} \mathfrak{g}_{\alpha}$$

Parabolic subgroups

$$\mathfrak{p} = \mathfrak{h} \oplus \bigoplus_{\alpha \in \Delta(\mathfrak{p})} \mathfrak{g}_{\alpha} \quad \Delta(\mathfrak{p}) = \Delta_+ \cup \langle \Sigma \rangle$$

Levi decomposition

$$\mathfrak{p} = \mathfrak{l} \oplus \mathfrak{u} \quad \mathfrak{l} = \mathfrak{h} \oplus \bigoplus_{\alpha \in \langle \Sigma \rangle} \mathfrak{g}_{\alpha} \quad \mathfrak{u} = \bigoplus_{\alpha \in \Delta(\mathfrak{u})} \mathfrak{g}_{\alpha}$$

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Parabolic subgroups

$$\mathfrak{p} = \mathfrak{l} \oplus \mathfrak{u} \quad \mathfrak{l} = \mathfrak{h} \oplus \bigoplus_{\alpha \in \langle \Sigma \rangle} \mathfrak{g}_\alpha \quad \mathfrak{u} = \bigoplus_{\alpha \in \Delta(\mathfrak{u})} \mathfrak{g}_\alpha$$

Corresponding group P

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Corresponding group P



Minimal parabolic
Borel

$$B = NA$$

$$N = \left\{ \begin{pmatrix} \boxed{1} & * & * & * \\ & \boxed{1} & * & * \\ & & \boxed{1} & * \\ & & & \boxed{1} \end{pmatrix} \right\}$$

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$$U^{(1)} = U \quad U^{(n+1)} = [U^{(n)}, U^{(n)}]$$

Terminology

$P = B \rightarrow U = N$ Fourier coefficient is a Whittaker coefficient

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F_U

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Iwasawa decomposition

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↑ Iwasawa decomposition

Characters and coefficients with all $m_\alpha \neq 0$ are called **generic**
otherwise they are called **degenerate**

Fourier expansion

Choice of unipotent subgroup U \longleftrightarrow Study different perturbative and non-perturbative effects

[Green-Miller-Vanhove]

Fourier expansion

Choice of unipotent subgroup U \longleftrightarrow Study different perturbative and non-perturbative effects

- String perturbation limit
D-instantons | NS5-instantons

$$g_s \rightarrow 0$$



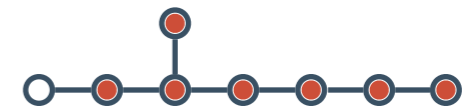
[Green-Miller-Vanhove]

Fourier expansion

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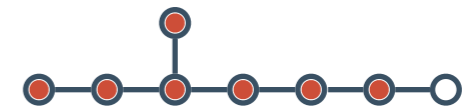
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- Decompactification limit
Higher dimensional black holes | BPS states

Large radius for
compactified circle



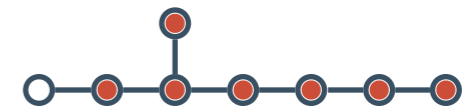
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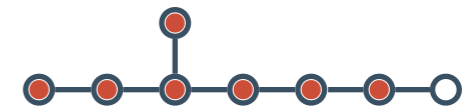
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- M-theory limit
M2, M5-instantons

Large M-theory torus



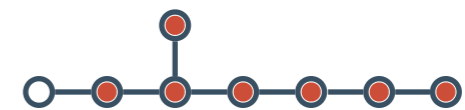
[Green-Miller-Vanhove]

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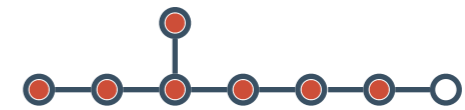
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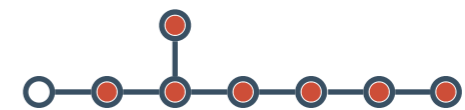
Maximal parabolic
subgroups

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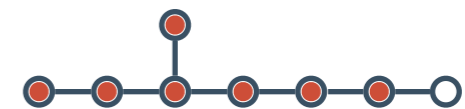
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[Green-Miller-Vanhove]

Maximal parabolic subgroups

Difficult to compute!

Fourier expansion

Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients

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Adelic framework

*An **efficient**, but abstract, way to approach the subject of automorphic forms is by the introduction of **adeles**, rather **ungainly objects** that nevertheless, once familiar, **spare** much unnecessary thought and **many useless calculations**.*

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Lift to the adèles

[FGKP15 §4.2.2]

$$\mathbb{A} = \mathbb{A}_{\mathbb{Q}} \quad G(\mathbb{A}) = G(\mathbb{R}) \times \prod'_{p \text{ prime}} G(\mathbb{Q}_p) \quad K_{\mathbb{A}} = K \times \prod_{p \text{ prime}} G(\mathbb{Z}_p)$$

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Eisenstein series \longrightarrow Adelic Eisenstein series

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Fourier coefficients \longrightarrow Adelic Fourier coefficients

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Fourier coefficients \longrightarrow Adelic Fourier coefficients

$$\int_{U(\mathbb{Z}) \backslash U(\mathbb{R})} E(\chi; ug) \overline{\psi_{\mathbb{R}}(u)} du$$

$$\int_{U(\mathbb{Q}) \backslash U(\mathbb{A})} E(\chi; ug) \overline{\psi_{\mathbb{A}}(u)} du$$

$$m_{\alpha} \in \mathbb{Z}$$

$$m_{\alpha} \in \mathbb{Q}$$

Computing adelic Fourier coefficients

[FGKP15 §9-10]

Whittaker coefficients

Computing adelic Fourier coefficients

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Constant term: Langlands' constant term formula

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Computing adelic Fourier coefficients

[FGKP15 §9-10]

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Generic coefficient: Factorises over the primes. Casselman-Shalika formula

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[GKP14]

Fourier coefficients

In terms of Whittaker coefficients

Simplify drastically for certain χ

Example of simplifications

$$G = SL(3)$$

$$E(\chi; g)$$

$$\chi \longleftrightarrow (s_1, s_2) \in \mathbb{C}^2$$

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$$\psi_N \left(\begin{pmatrix} 1 & x_1 & * \\ & 1 & x_2 \\ & & 1 \end{pmatrix} \right) = e^{2\pi i(m_1 x_1 + m_2 x_2)}$$

$$\begin{matrix} m_1 & m_2 \\ \circ & \text{---} & \circ \end{matrix}$$

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$$W_N(\chi, \psi_{m_1, m_2}; g) \propto \left(\begin{matrix} \text{arithmetic} \\ \text{factor} \end{matrix} \right) \int K_{\#}(\dots) K_{\#}(\dots)$$

[FGKP15 §10.6]

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p-adic part


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Vanishes for certain (s_1, s_2)

[FGKP15 §10.6]

Example of simplifications

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Example of simplifications

Certain (s_1, s_2)

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$$W_N(\chi, \psi_{m_1, 0}; g) \propto K_{\#}(\dots) + K_{\#}(\dots) + K_{\#}(\dots)$$

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Automorphic representations

$G(\mathbb{A}) \curvearrowright$ Space of automorphic forms*

* With some subtleties described in [FGKP15 §6]

[Bump, Goldfeld-Hundley]

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$G(\mathbb{A}) \curvearrowright$ Space of automorphic forms*

Automorphic representation π = an irreducible component of the above space under this action

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Automorphic representations

$G(\mathbb{A}) \curvearrowright$ Space of automorphic forms*

Automorphic representation π = an irreducible component of the above space under this action

What is a small automorphic representation?

* With some subtleties described in [FGKP15 §6]

[Bump, Goldfeld-Hundley]

Wavefront set

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg,
Gomez-Gourevitch-Sahi, Jiang-Liu-Savin, Joseph, Miller-Sahi]

Wavefront set

The (global) wavefront set contains all the characters ψ which can give rise to non-vanishing Fourier coefficients in that representation

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg, Gomez-Gourevitch-Sahi, Jiang-Liu-Savin, Joseph, Miller-Sahi]

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Small automorphic representations have few non-vanishing Fourier coefficients

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg, Gomez-Gourevitch-Sahi, Jiang-Liu-Savin, Joseph, Miller-Sahi]

Wavefront set

Characters ψ \longleftrightarrow Nilpotent elements in \mathfrak{g}

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg,
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Wavefront set

Characters $\psi \longleftrightarrow$ Nilpotent elements in \mathfrak{g}

Nilpotent orbit $\mathcal{O} = \{gXg^{-1} \mid g \in G(\mathbb{C})\}$ $X \in \mathfrak{g}$ nilpotent

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg,
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↑
So called special orbits

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg, Gomez-Gourevitch-Sahi, Jiang-Liu-Savin, Joseph, Miller-Sahi]

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Closure with respect to partial ordering

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Nilpotent orbits

[Collingwood-McGovern]

For $SL(n)$, orbits can be identified
with partitions of n

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(p_1, p_2, \dots)

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For $SL(n)$, orbits can be identified with partitions of n

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 $(p_1, p_2, \dots) \leq (q_1, q_2, \dots)$ partial ordering

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$(p_1, p_2, \dots) \leq (q_1, q_2, \dots)$ partial ordering

\iff

$$\sum_{i=1}^k p_i \leq \sum_{i=1}^k q_i \quad \forall k$$

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Illustrated by a Hasse diagram

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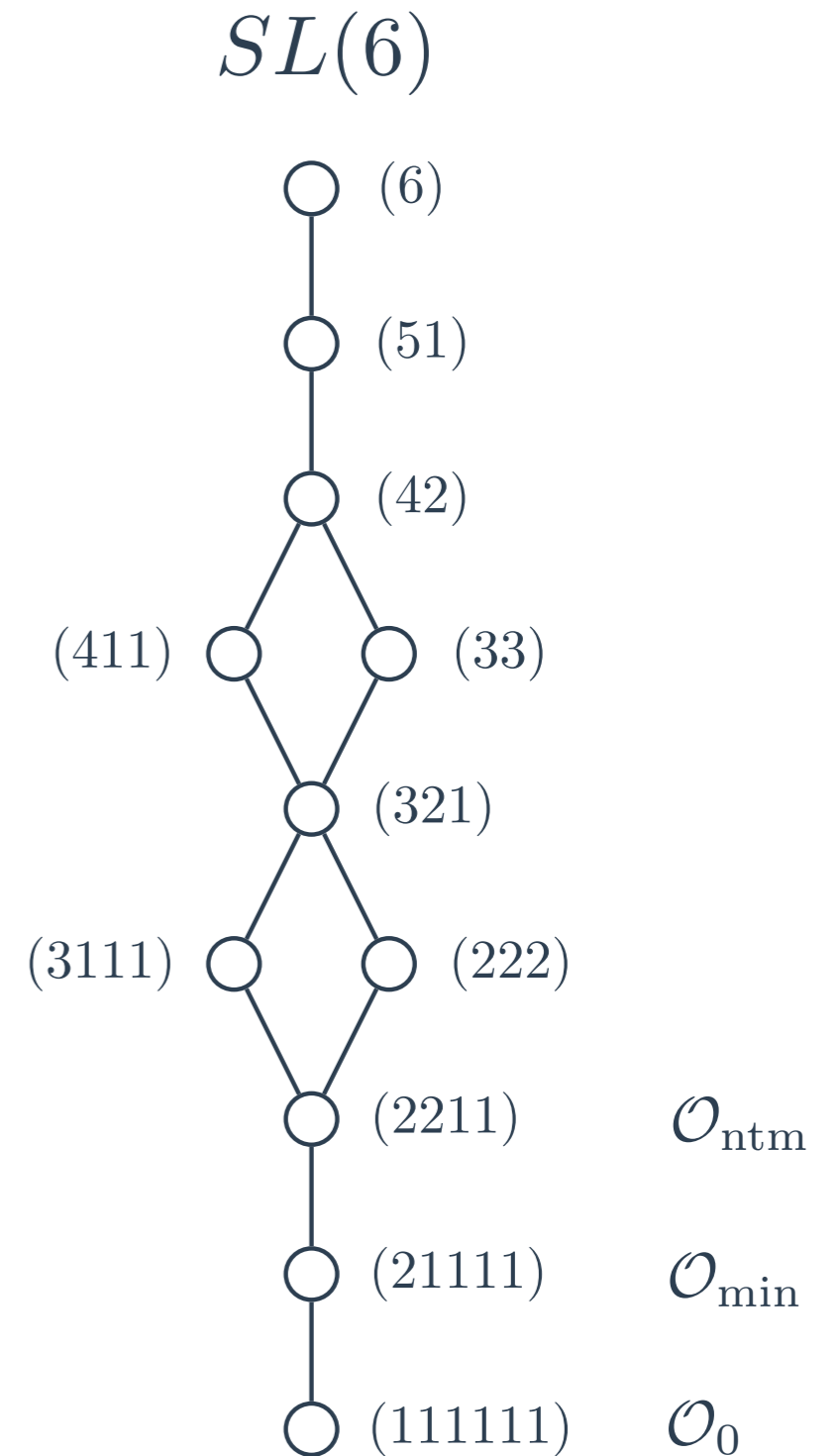
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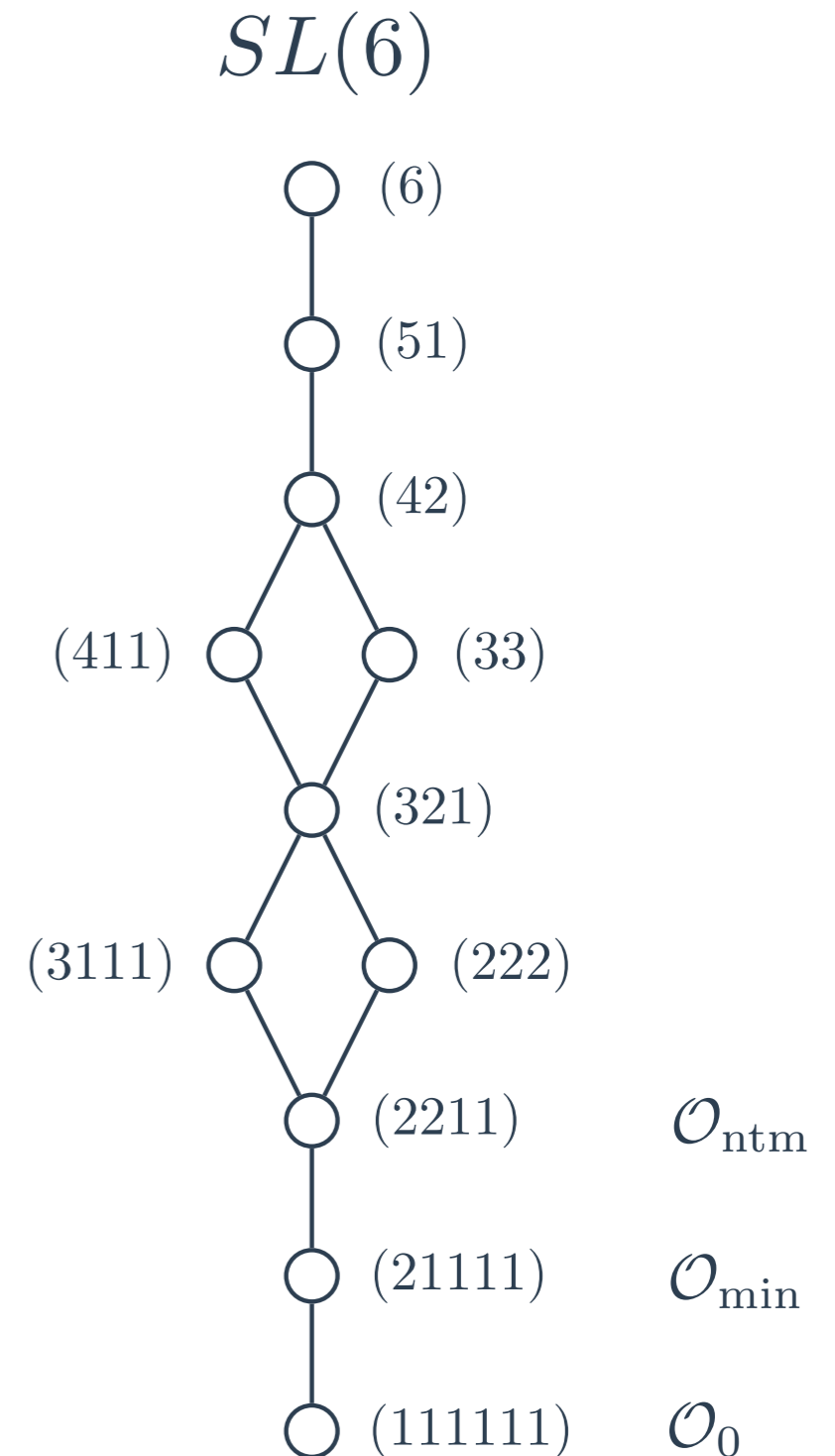
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Illustrated by a Hasse diagram

Closure: $\overline{\mathcal{O}} = \bigcup_{\mathcal{O}' \leq \mathcal{O}} \mathcal{O}'$



Automorphic representations

Small representations

Automorphic representations

Small representations

$$\mathrm{WF}(\pi_{\min}) = \overline{\mathcal{O}_{\min}} = \mathcal{O}_{\min} \cup \mathcal{O}_0$$

$$\mathrm{WF}(\pi_{\mathrm{ntm}}) = \overline{\mathcal{O}_{\mathrm{ntm}}} = \mathcal{O}_{\mathrm{ntm}} \cup \mathcal{O}_{\min} \cup \mathcal{O}_0$$

Automorphic representations

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$$\mathcal{E}_{(0,0)}^{(D)} \in \pi_{\min}$$

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[Green-Miller-Vanhove,
Pioline, Bossard-Verschinin]

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Certain $(s_1, s_2) \longleftrightarrow \chi_{\min}$ such that $E(\chi_{\min}, g) \in \pi_{\min}$

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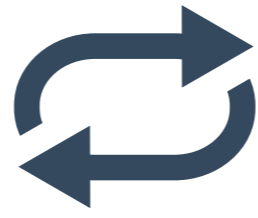
$$\mathcal{E}_{(1,0)}^{(D)} \in \pi_{\mathrm{ntm}}$$

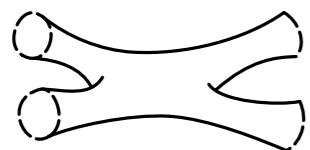
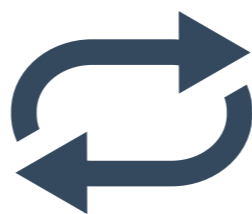
[Green-Miller-Vanhove,
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Certain $(s_1, s_2) \longleftrightarrow \chi_{\min}$ such that $E(\chi_{\min}, g) \in \pi_{\min}$

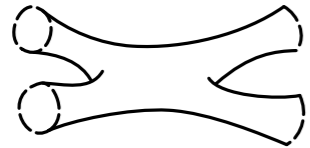
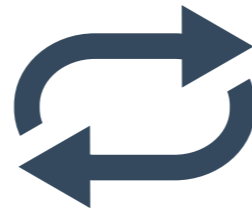
$$\int_K K \longrightarrow 0$$

$$\sum_K K \longrightarrow K$$



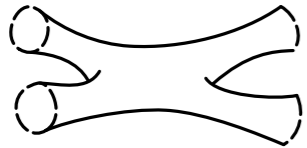
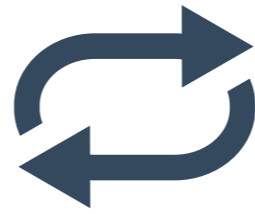


$$R + (\alpha')^3 \mathcal{E}_{(0,0)}^{(D)}(g) R^4 + (\alpha')^5 \mathcal{E}_{(1,0)}^{(D)}(g) D^4 R^4 + (\alpha')^6 \mathcal{E}_{(0,1)}^{(D)}(g) D^6 R^4 + \dots$$



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D	$G(\mathbb{R})$	K	$G(\mathbb{Z})$
10	$SL(2, \mathbb{R})$	$SO(2)$	$SL(2, \mathbb{Z})$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$	$SL(2, \mathbb{Z}) \times \mathbb{Z}_2$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$	$SL(3, \mathbb{Z}) \times SL(2, \mathbb{Z})$
7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5)) / \mathbb{Z}_2$	$Spin(5, 5; \mathbb{Z})$
5	$E_6(\mathbb{R})$	$USp(8) / \mathbb{Z}_2$	$E_6(\mathbb{Z})$
4	$E_7(\mathbb{R})$	$SU(8) / \mathbb{Z}_2$	$E_7(\mathbb{Z})$
3	$E_8(\mathbb{R})$	$Spin(16) / \mathbb{Z}_2$	$E_8(\mathbb{Z})$



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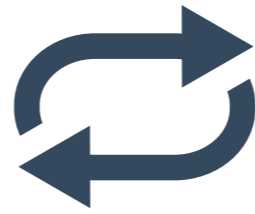
$$\mathcal{E}_{(0,0)}(\tau) = 2\zeta(3)E(3/2; \tau)$$

$$\mathcal{E}_{(1,0)}(\tau) = \zeta(5)E(5/2; \tau)$$

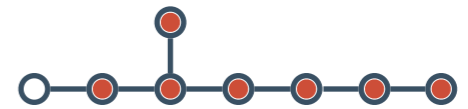




$$F_U(\chi, \psi; g) = \int_{U(\mathbb{Q}) \setminus U(\mathbb{A})} E(\chi; ug) \overline{\psi(u)} du$$



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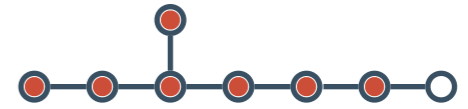


Maximal parabolic
subgroups

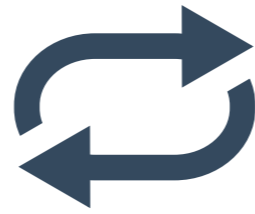


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$$\mathbb{Z} \backslash \mathbb{R} \longrightarrow \mathbb{Q} \backslash \mathbb{A}$$



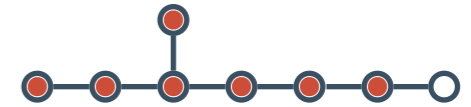
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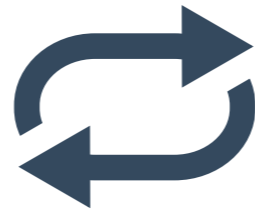
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Known Whittaker coefficients W_N



Maximal parabolic
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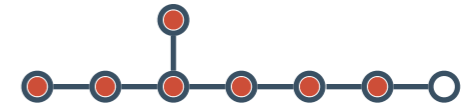
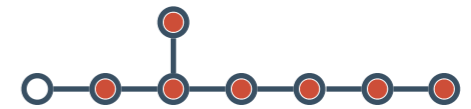


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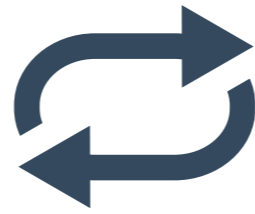
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Known Whittaker coefficients W_N

Automorphic representation π



Maximal parabolic
subgroups



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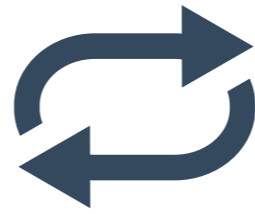
Known Whittaker coefficients W_N

Automorphic representation π

Vanishing properties $\text{WF}(\pi)$



Maximal parabolic
subgroups



Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients
using vanishing properties of the given π

Previous results

[Miller-Sahi]

Previous results

Theorem

For $G = E_6, E_7$, an automorphic form $\varphi \in \pi_{\min}$ is completely determined by maximally degenerate Whittaker coefficients

W_N with only one $m_\alpha \neq 0$

[Miller-Sahi]

Main results

SL(3), SL(4)

[GKP14]

Main results

$SL(3), SL(4)$

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[GKP14]

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More generally, for $\varphi \in \pi$

[GKP14]

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$$\varphi = \sum_{\mathcal{O}} \varphi_{\mathcal{O}} \quad \text{where } \varphi_{\mathcal{O}} \text{ vanishes unless } \mathcal{O} \subseteq \text{WF}(\pi)$$

[GKP14]

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$SL(3), SL(4)$

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Corollary

[GKP14]

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$SL(3), SL(4)$

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Corollary

$\varphi \in \pi_{\min}$ maximally degenerate Whittaker coefficients

[GKP14]

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Corollary

$\varphi \in \pi_{\min}$ maximally degenerate Whittaker coefficients single root

[GKP14]

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[GKP14]

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Corollary

$$\varphi \in \pi_{\min}$$

single root

$$\varphi \in \pi_{\text{ntm}}$$

at most two commuting roots

[GKP14]

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at most two commuting roots

.....
strongly orthogonal

[GKP14]

Main results

$SL(3), SL(4)$

Fourier coefficients on maximal parabolic subgroups



Main results

$SL(3), SL(4)$

Fourier coefficients on maximal parabolic subgroups



in the minimal representation

Main results

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in the minimal representation

$$\mathcal{E}_{(0,0)}^{(D)} \in \pi_{\min}$$

Main results

$SL(3), SL(4)$



$$\mathcal{E}_{(0,0)}^{(D)} \in \pi_{\min}$$

[GKP14]

Main results

$SL(3), SL(4)$



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$$F_U(\chi_{\min}, \psi; g) = W_N(\chi_{\min}, \psi'; lg) \quad \text{with } l \in L(\mathbb{Q}) \text{ depending on } \psi$$

[GKP14]

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$SL(3), SL(4)$



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↑
Maximal parabolic
Fourier coefficient

[GKP14]

Main results

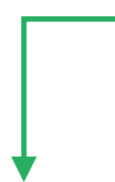
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Maximal parabolic
Fourier coefficient



[GKP14]

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Maximal parabolic
Fourier coefficient



Maximally degenerate

[GKP14]

Other groups

[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

Other groups

$$SL(n)$$

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Other groups

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↑
Maximal parabolic
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↑
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[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

Other groups

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↑
Maximal parabolic
Fourier coefficient

↑
Maximally degenerate

and similar statement for next-to-minimal representation

[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

Other groups

Conjecture

A similar relations holds for all simple, simply laced Lie groups

$$F_U(\chi_{\min}, \psi; g) = W_N(\chi_{\min}, \psi'; lg) \quad \text{with } l \in L(\mathbb{Q}) \text{ depending on } \psi$$

↑
Maximal parabolic
Fourier coefficient

↑
Maximally degenerate

[GKP14]

[Proof in progress with Gourevitch, Kleinschmidt, Persson, Sahi]

Local spherical vectors

Checks for E_6, E_7, E_8

Local spherical vectors

Checks for E_6, E_7, E_8

$$\pi_{\min,p} \subset \text{Ind}_{P(\mathbb{Q}_p)}^{G(\mathbb{Q}_p)} \chi_{\min,p} \hookrightarrow \text{Ind}_{U(\mathbb{Q}_p)}^{G(\mathbb{Q}_p)} \psi_{U,p}$$

multiplicity one

[Gan-Savin]

Local spherical vectors

Checks for E_6, E_7, E_8

$$\pi_{\min,p} \subset \text{Ind}_{P(\mathbb{Q}_p)}^{G(\mathbb{Q}_p)} \chi_{\min,p} \hookrightarrow \text{Ind}_{U(\mathbb{Q}_p)}^{G(\mathbb{Q}_p)} \psi_{U,p} \quad \text{multiplicity one}$$

[Gan-Savin]

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$$f_{\psi_{U,p}}^\circ \in \text{Ind}_{U(\mathbb{Q}_p)}^{G(\mathbb{Q}_p)} \psi_{U,p} \quad \text{computed in several cases} \quad p \leq \infty$$

[Dvorsky-Sahi, Kazhdan-Polishchuk, Kazhdan-Pioline, Savin-Woodbury]

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Complete **agreement** for E_6, E_7, E_8 in both abelian and Heisenberg realisations

Outlook

Tools for proving the conjecture

[Work in progress with Gourevitch, Kleinschmidt, Persson, Sahi]

Outlook

Tools for proving the conjecture

$(S, \psi) \in \mathfrak{g} \times \mathfrak{g}^*$ Whittaker pair

[Gomez-Gourevitch-Sahj]

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— Semi-simple

Whittaker pair

[Gomez-Gourevitch-Sahj]

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Methods for relating different Whittaker pairs

$$(S, \psi) \longrightarrow (S', \psi')$$

[Work in progress with Gourevitch, Kleinschmidt, Persson, Sahj]

Outlook



Outlook

Simplification of Fourier coefficients with χ_{\min} for dimensions lower than three. Kac-Moody groups E_9, E_{10}, E_{11}

[Fleig-Kleinschmidt, Fleig-Kleinschmidt-Persson]

How to define “small automorphic representations” for Kac-Moody groups? What is the mechanism behind the vanishing properties?



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[Fleig-Kleinschmidt, Fleig-Kleinschmidt-Persson]

How to define “small automorphic representations” for Kac-Moody groups? What is the mechanism behind the vanishing properties?

$\mathcal{E}_{(0,1)} D^6 R^4$ requires extended notion of automorphic forms, the development of which will positively bring new exciting insights to both physics and mathematics.



Thank you!

Henrik Gustafsson

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